

Today's summary

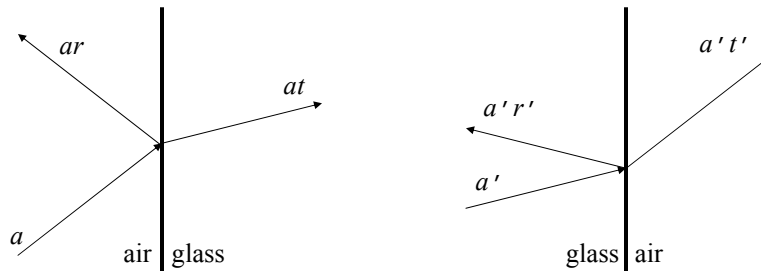
- Multiple beam interferometers: Fabry-Perot resonators
 - Stokes relationships
 - Transmission and reflection coefficients for a dielectric slab
 - Optical resonance
- Principles of lasers
- Coherence: spatial / temporal

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Fabry-Perot interferometers

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Relation between r, r' and t, t'



$$r' = -r$$

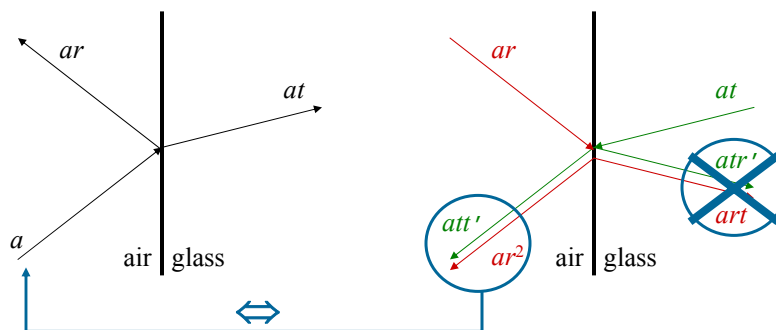
$$r^2 + tt' = 1$$

Stokes relationships

Proof: algebraic from the Fresnel coefficients
or using the property of *preservation of the field properties upon time reversal*

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Proof using time reversal

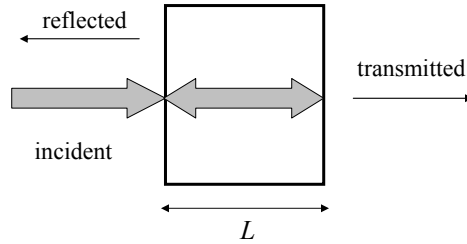


$$a(r + r')t = 0 \Rightarrow r = -r'$$

$$a(r^2 + tt') = a \Rightarrow r^2 + tt' = 1$$

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Fabry-Perot Interferometer



Resonance condition: reflected wave = 0

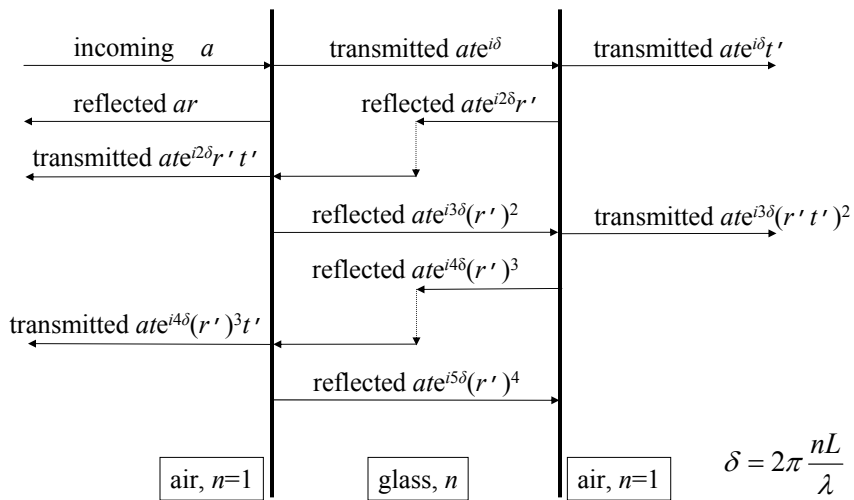
⇔ all reflected waves interfere destructively

$$L = \frac{m\lambda}{2n}$$

← wavelength in free space
← refractive index

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Calculation of the reflected wave



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Calculation of the reflected wave

$$\begin{aligned} a_{\text{reflected}} &= a \left\{ r + tt' r' e^{i2\delta} \left(1 + r'^2 e^{i2\delta} + r'^4 e^{i4\delta} + \dots \right) \right\} \\ &= a \left\{ r + tt' r' e^{i2\delta} \frac{1}{1 - r'^2 e^{i2\delta}} \right\} \end{aligned}$$

Use Stokes relationships $r' = -r$
 $r^2 + tt' = 1$

$$a_{\text{reflected}} = a \frac{r(1 - e^{i2\delta})}{1 - r^2 e^{i2\delta}}$$

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Transmission & reflection coefficients

$$a_{\text{reflected}} = a \frac{(1 - e^{i2\delta})r}{1 - r^2 e^{i2\delta}} \quad a_{\text{transmitted}} = a \frac{tt'}{1 - r^2 e^{i2\delta}}$$

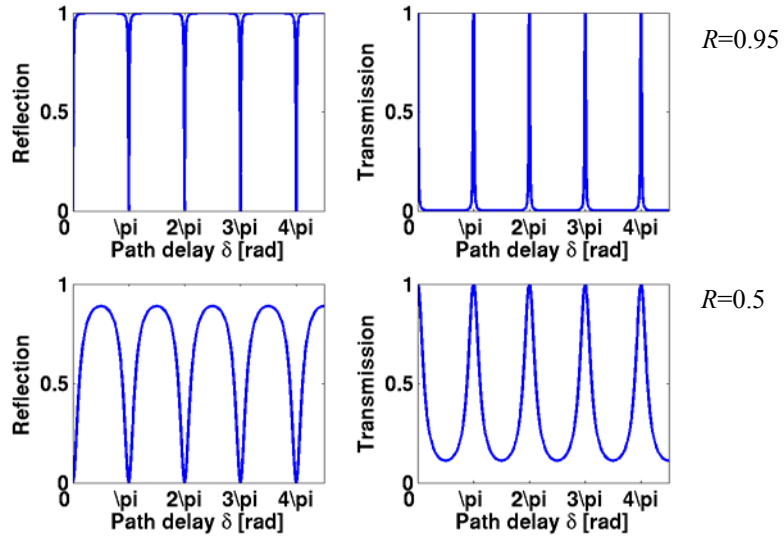
$$R \equiv |r|^2$$

$$\left(\begin{array}{l} \text{reflection} \\ \text{coefficient} \end{array} \right) = \left| \frac{a_{\text{reflected}}}{a} \right|^2 = \frac{4R \sin^2 \delta}{(1 - R)^2 + 4R \sin^2 \delta}$$

$$\left(\begin{array}{l} \text{transmission} \\ \text{coefficient} \end{array} \right) = \left| \frac{a_{\text{transmitted}}}{a} \right|^2 = \frac{(1 - R)^2}{(1 - R)^2 + 4R \sin^2 \delta}$$

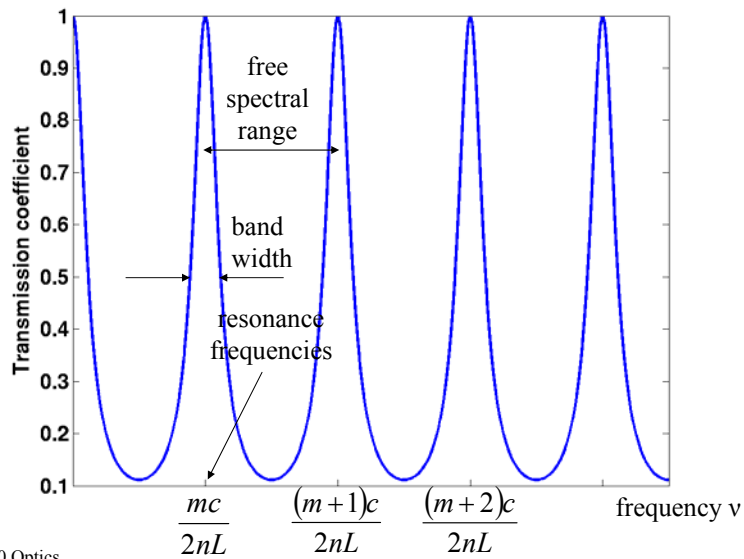
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Transmission & reflection vs path



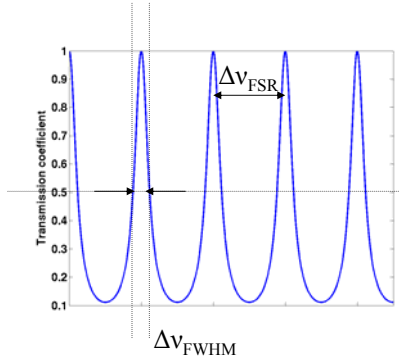
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Fabry-Perot terminology



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Fabry-Perot terminology



FWHM Bandwidth is inversely proportional to the *finesse* F (or *quality factor*) of the cavity

$$F \equiv \frac{\pi\sqrt{R}}{1-R}$$

$$\Delta\nu_{\text{FWHM}} = \frac{c}{2nLF}$$

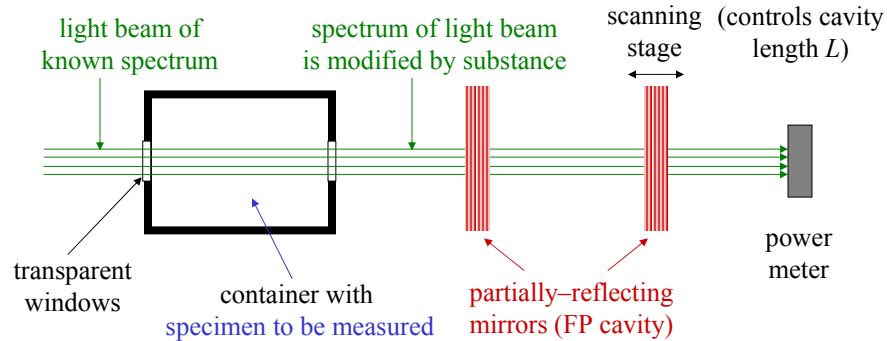
$$\Delta\nu_{\text{FWHM}} = \frac{\Delta\nu_{\text{FSR}}}{F} \quad (\text{bandwidth}) = \frac{(\text{free spectral range})}{(\text{finesse})}$$

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Spectroscopy using Fabry-Perot cavity

Goal: to measure the specimen's absorption as function of frequency ω

Experimental measurement principle:

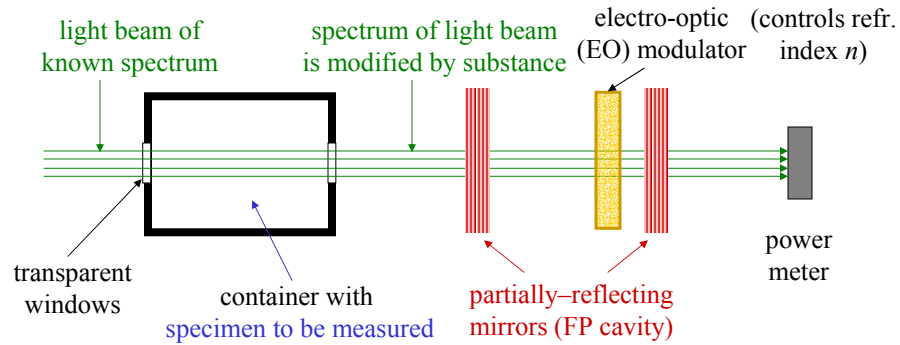


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Spectroscopy using Fabry-Perot cavity

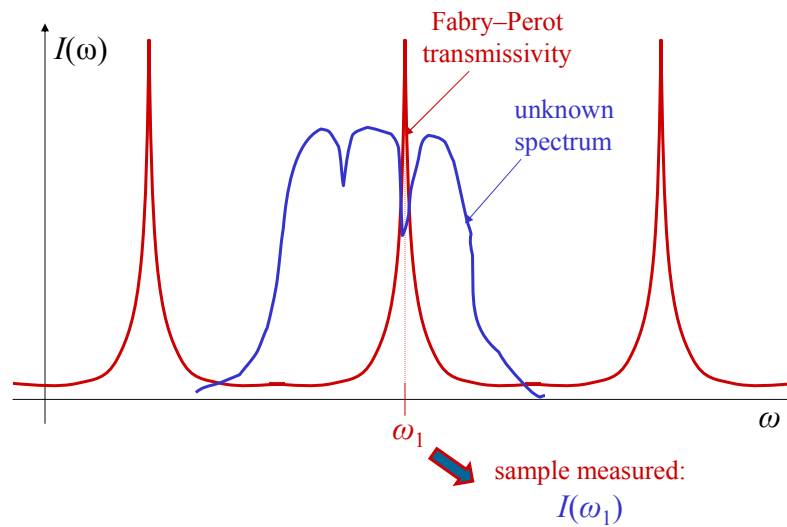
Goal: to measure the specimen's absorption as function of frequency ω

Experimental measurement principle:



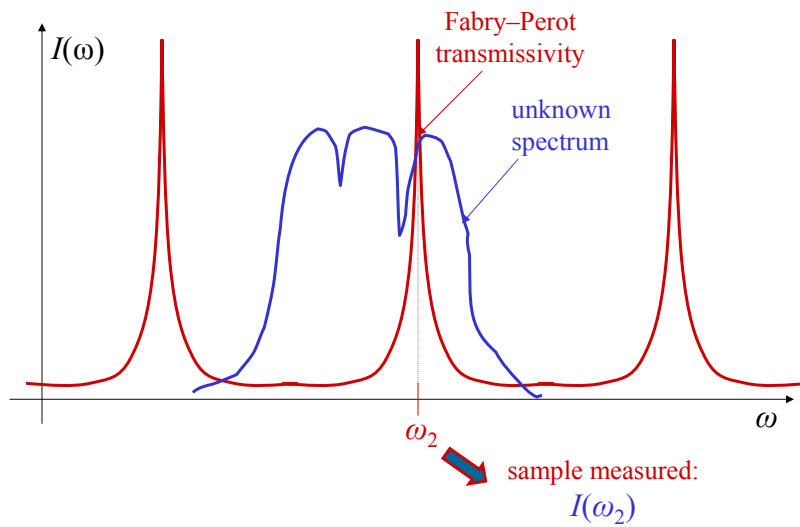
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Spectroscopy using Fabry-Perot cavity



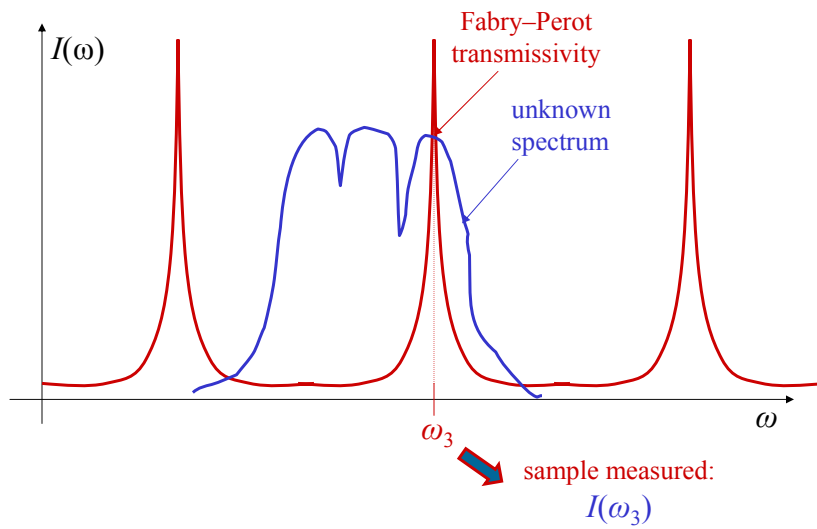
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Spectroscopy using Fabry-Perot cavity



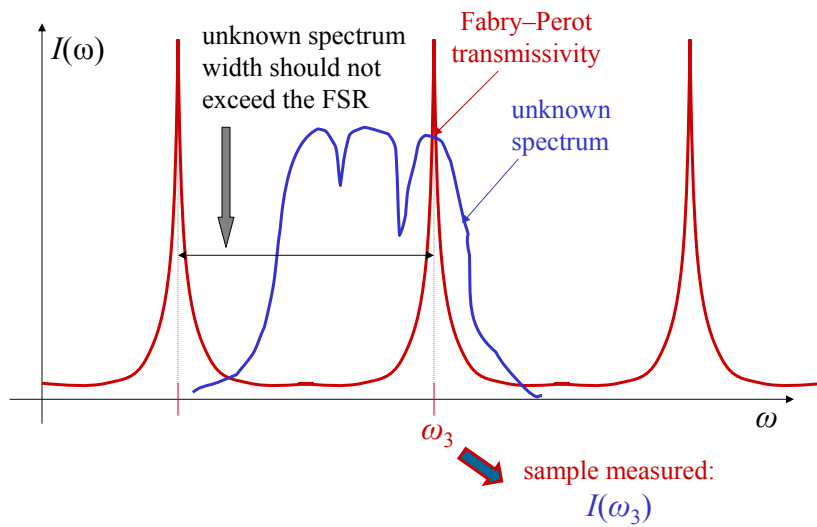
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Spectroscopy using Fabry-Perot cavity



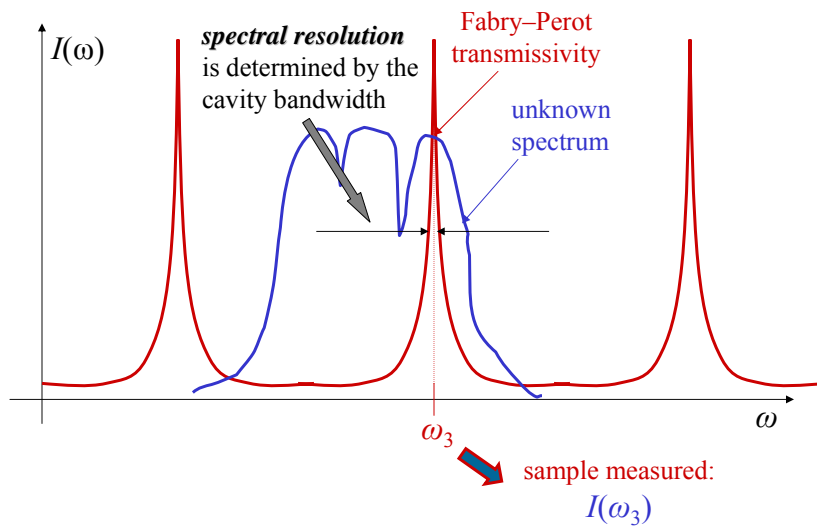
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Spectroscopy using Fabry-Perot cavity



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Spectroscopy using Fabry-Perot cavity

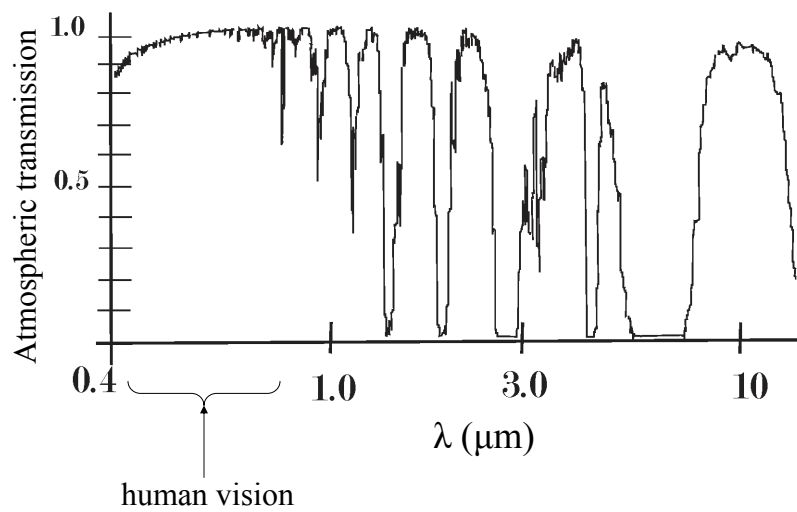


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Lasers

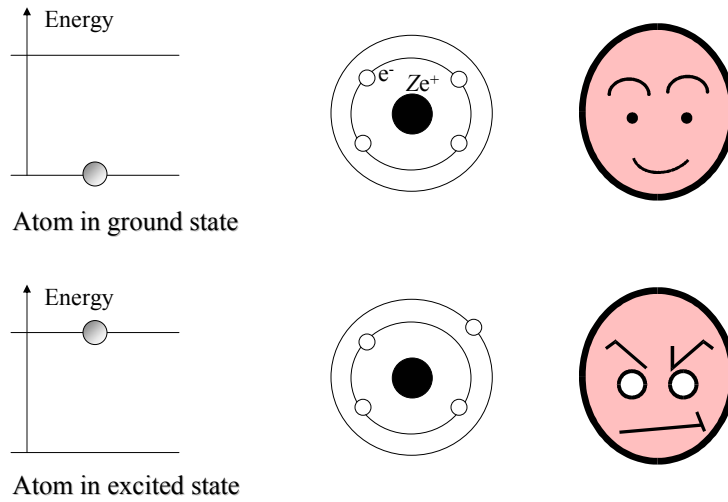
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Absorption spectra



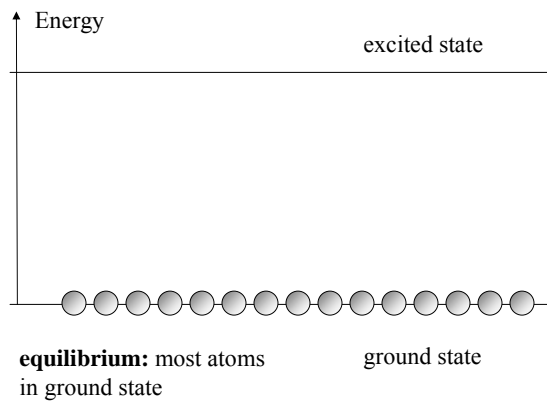
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Semi-classical view of atom excitations



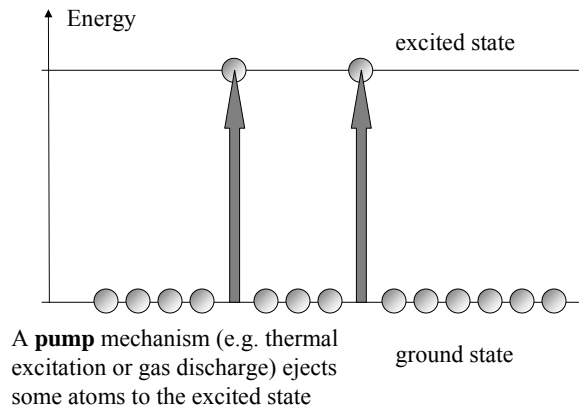
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Light generation



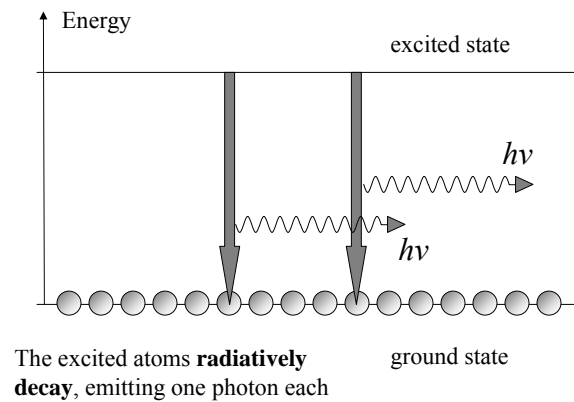
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Light generation



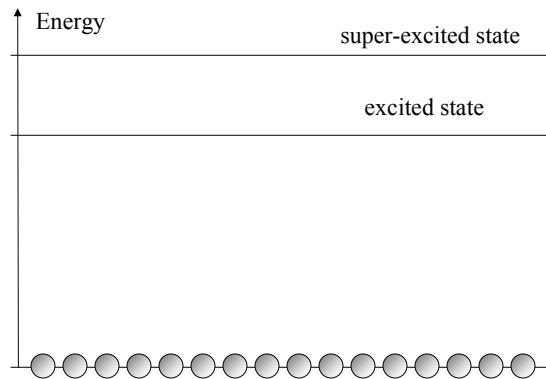
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Light generation



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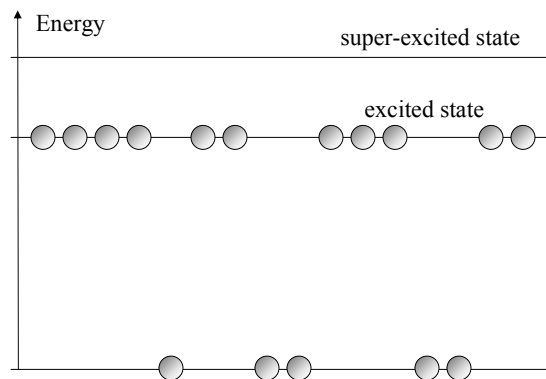
Light amplification: 3-level system



equilibrium: most atoms in ground state; note the existence of a third, "super-excited" state

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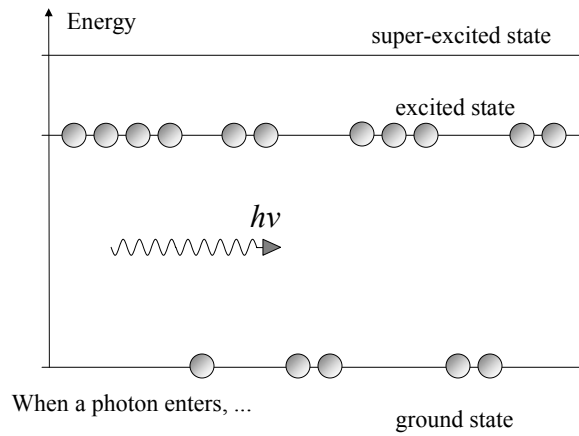
Light amplification: 3-level system



Utilizing the super-excited state as a short-lived "pivot point," the pump creates a **population inversion**

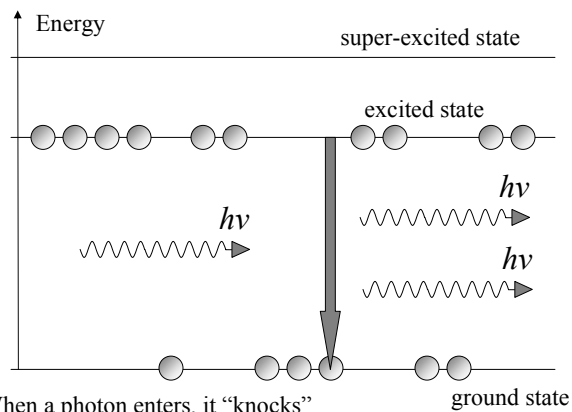
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Light amplification: 3-level system



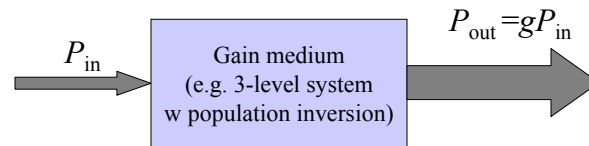
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Light amplification: 3-level system



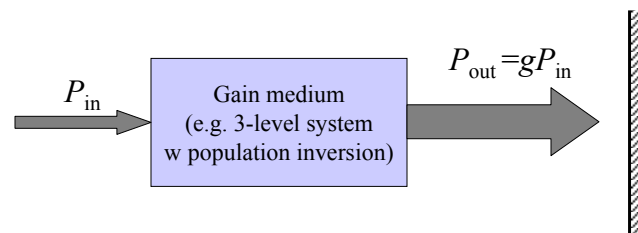
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Light amplifier

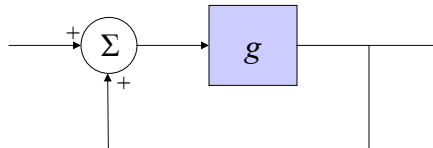


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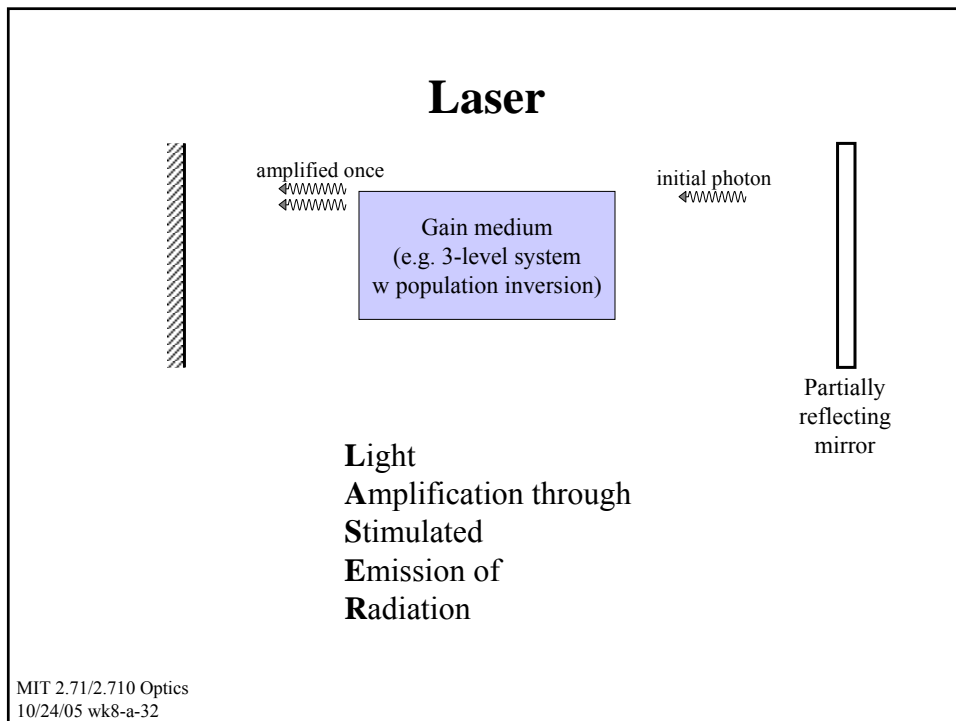
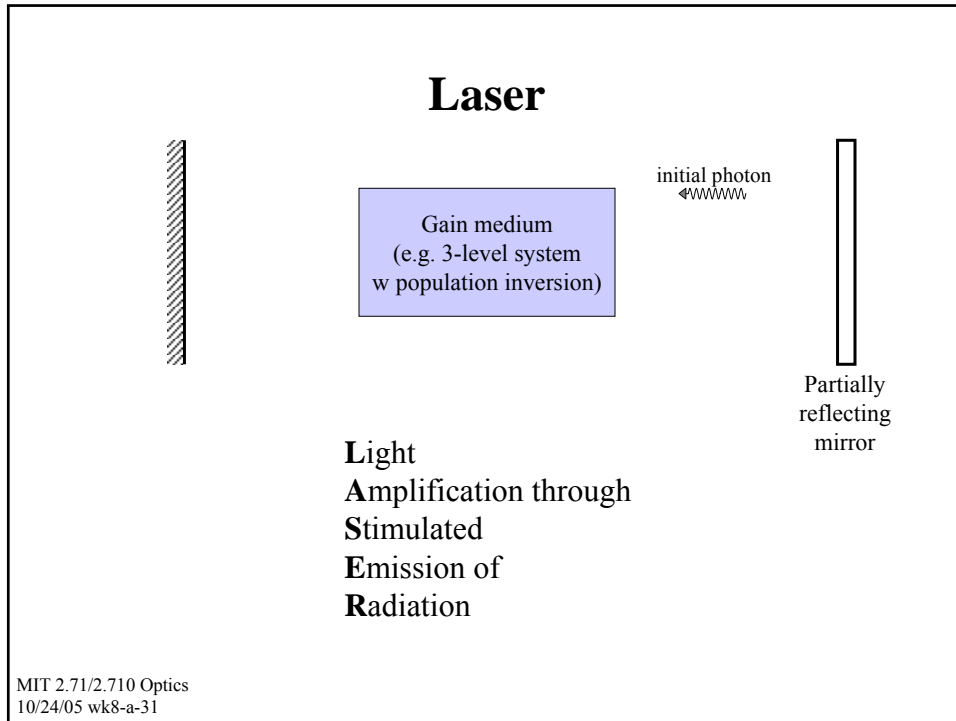
Light amplifier w positive feedback

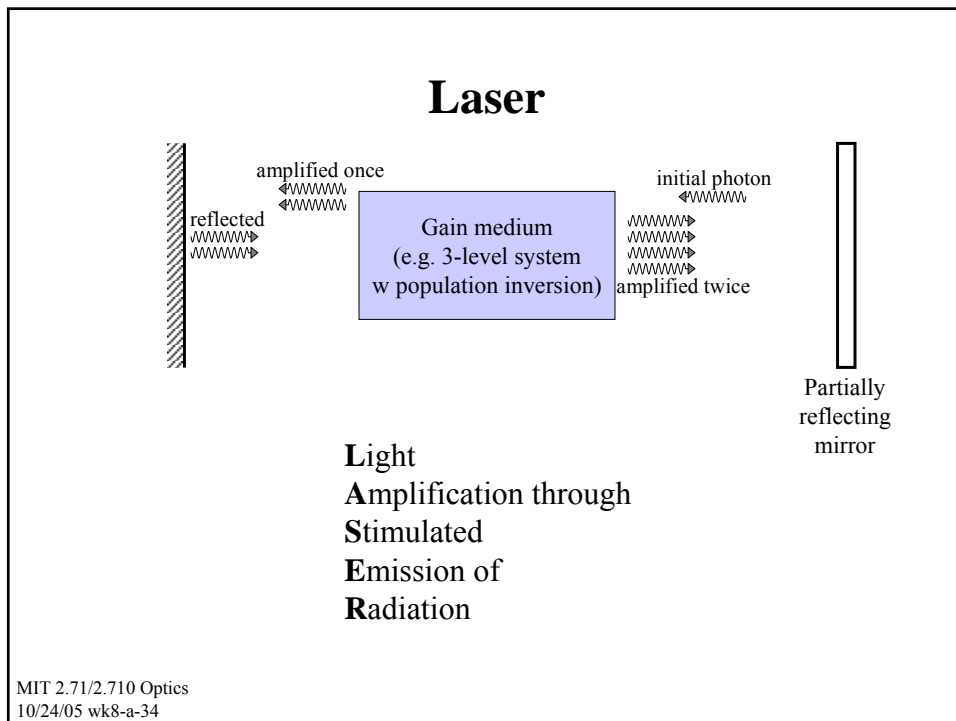
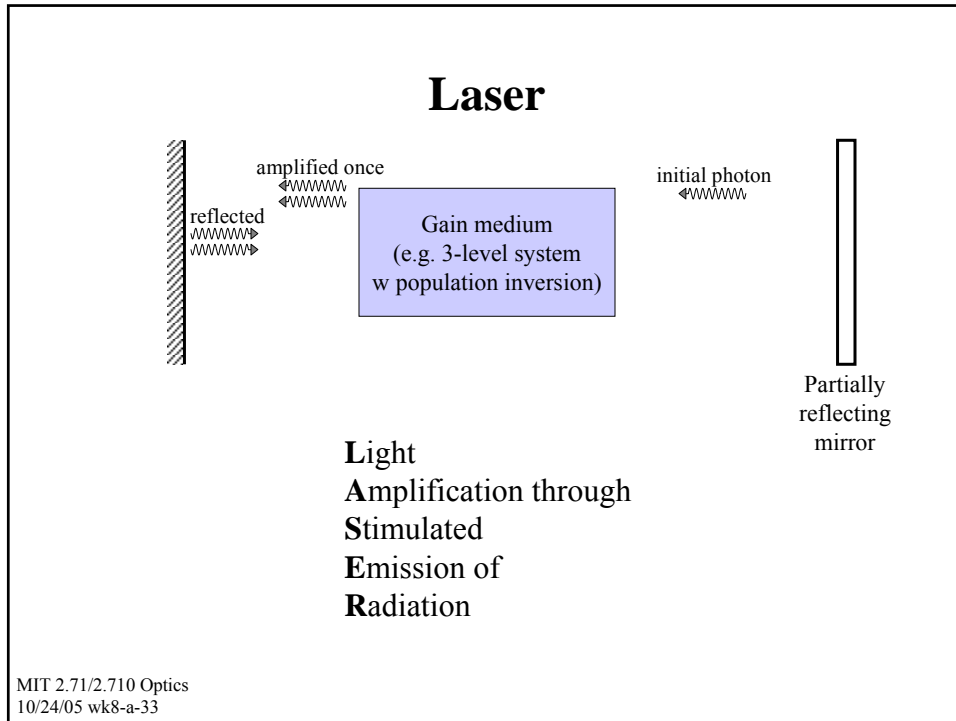


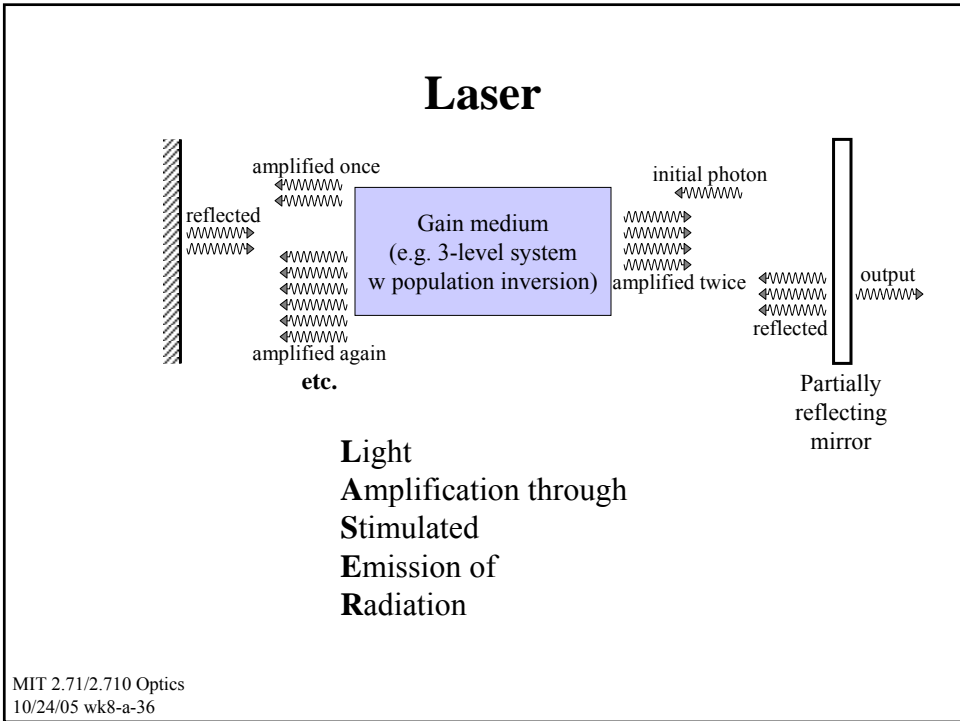
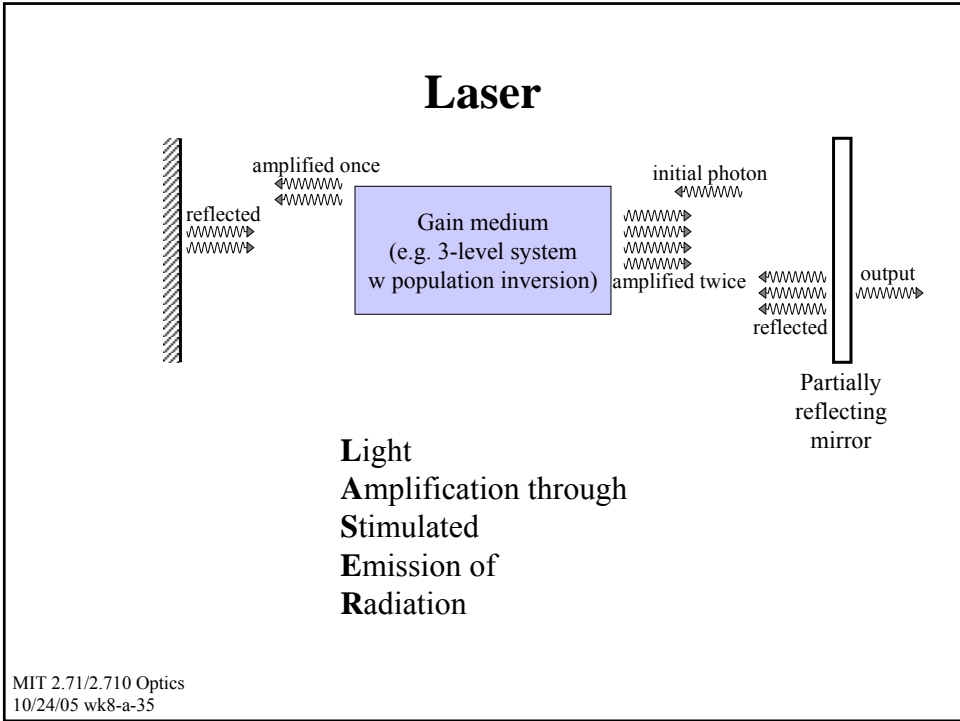
When the gain exceeds the roundtrip losses, the system goes into **oscillation**



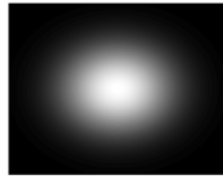
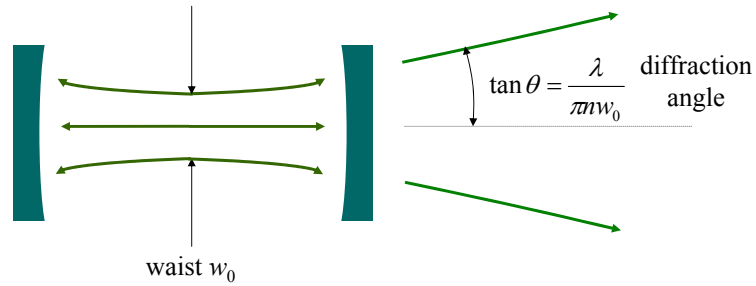
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Confocal laser cavities

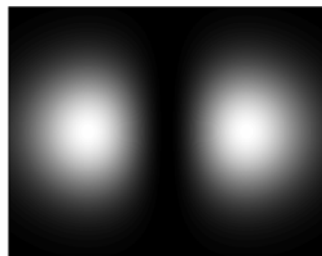


Beam profile:
2D Gaussian function

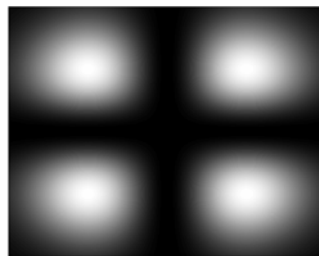
“TE₀₀ mode”

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Other “transverse modes”



TE10



TE11

(usually undesirable)

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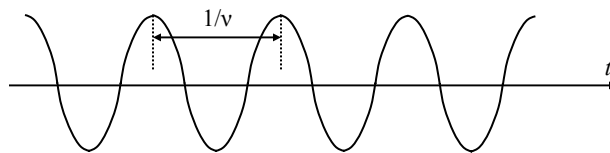
Types of lasers

- Continuous wave (cw)
- Pulsed
 - Q-switched
 - mode-locked
- Gas (Ar-ion, HeNe, CO₂)
- Solid state (Ruby, Nd:YAG, Ti:Sa)
- Diode (semiconductor)
- Vertical cavity surface-emitting lasers –VCSEL– (also sc)
- Excimer (usually ultra-violet)

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CW (continuous wave lasers)

Laser oscillation well approximated by a sinusoid



Typical sources:

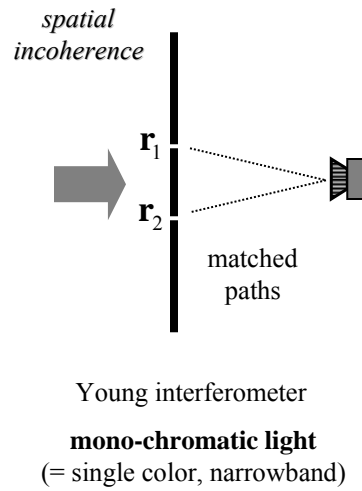
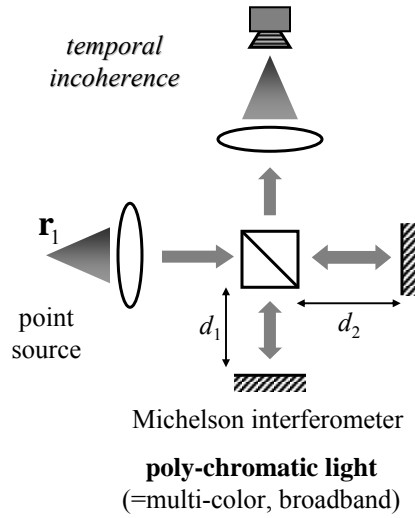
- Argon-ion: 488nm (blue) or 514nm (green); power ~1-20W
- Helium-Neon (HeNe): 633nm (red), also in green and yellow; ~1-100mW
- doubled Nd:YAG: 532nm (green); ~1-10W

Quality of sinusoid maintained over a time duration known as
“coherence time” t_c

Typical coherence times ~20nsec (HeNe), ~10μsec (doubled Nd:YAG)

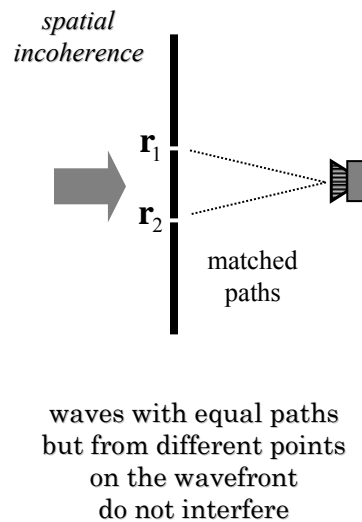
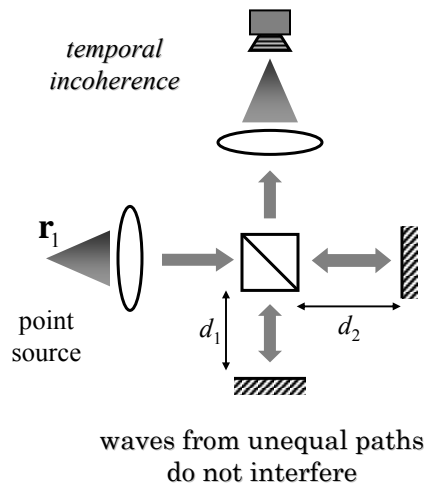
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Two types of incoherence



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Two types of incoherence



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Coherent vs incoherent beams

$$a_1 = |a_1|e^{i\phi_1}$$

$$a_2 = |a_2|e^{i\phi_2}$$

Mutually coherent: superposition field *amplitude* is described by *sum of complex amplitudes*

$$a = a_1 + a_2 = |a_1|e^{i\phi_1} + |a_2|e^{i\phi_2}$$

$$I = |a|^2 = |a_1 + a_2|^2$$

$$I_1$$

$$I_2$$

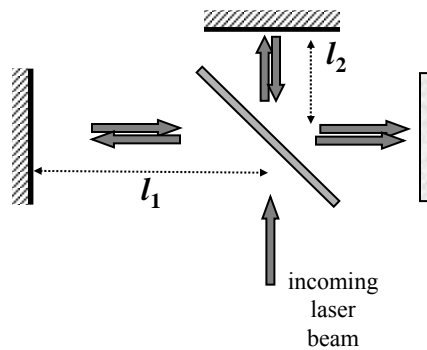
Mutually incoherent: superposition field *intensity* is described by *sum of intensities*

$$I = I_1 + I_2$$

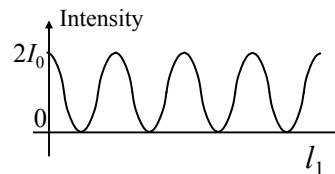
(the phases of the individual beams vary randomly with respect to each other; hence, we would need statistical formulation to describe them properly — *statistical optics*)

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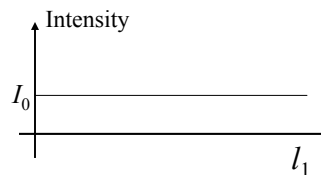
Coherence time and coherence length



- $l_1 - l_2$ much shorter than “coherence length” ct_c
sharp interference fringes

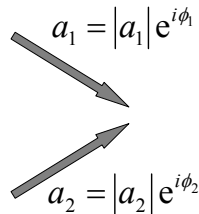


- $l_1 - l_2$ much longer than “coherence length” ct_c
no interference



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Coherent vs incoherent beams



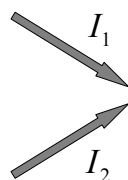
$$a_1 = |a_1| e^{i\phi_1}$$

$$a_2 = |a_2| e^{i\phi_2}$$

Coherent: superposition field *amplitude* is described by *sum of complex amplitudes*

$$a = a_1 + a_2 = |a_1| e^{i\phi_1} + |a_2| e^{i\phi_2}$$

$$I = |a|^2 = |a_1 + a_2|^2$$



$$I_1$$

$$I_2$$

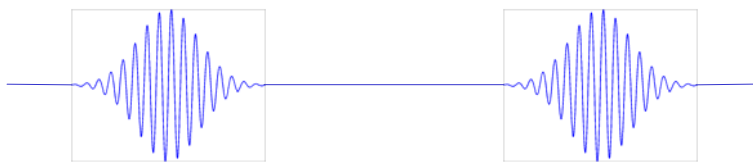
Incoherent: superposition field *intensity* is described by *sum of intensities*

$$I = I_1 + I_2$$

(the phases of the individual beams vary randomly with respect to each other; hence, we would need statistical formulation to describe them properly — *statistical optics*)

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Mode-locked lasers



Typical sources: Ti:Sa lasers (major vendors: Coherent, Spectra Phys.)

Typical mean wavelengths: 700nm – 1.4 μ m (near IR)

can be doubled to visible wavelengths

or split to visible + mid IR wavelengths using OPOs or OPAs

(OPO=optical parametric oscillator;

OPA=optical parametric amplifier)

Typical pulse durations: ~psec to few fsec

(just a few optical cycles)

Typical pulse repetition rates (“rep rates”): 80-100MHz

Typical average power: 1-2W; peak power ~MW-GW

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Overview of light sources

non-Laser

Thermal: polychromatic,
spatially incoherent
(e.g. light bulb)

Gas discharge: monochromatic,
spatially incoherent
(e.g. Na lamp)

Light emitting diodes (LEDs):
monochromatic, spatially
incoherent

Laser

Continuous wave (or cw):
strictly monochromatic,
spatially coherent
(e.g. HeNe, Ar⁺, laser diodes)

Pulsed: quasi-monochromatic,
spatially coherent
(e.g. Q-switched, mode-locked)

~nsec ~psec to few fsec
↑ ↑
pulse duration

mono/poly-chromatic = single/multi color